QUASI-3D MMF and Global CRM

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- I. Unification of Physics between GCM and CRM
	- 1. Introduction motivation and goal
	- 2. Strategy for development of a quasi-3D algorithm
	- 3. Current major issues
- II. Unification of Dynamics between GCM and CRM

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Results of preliminary experments and future plans

- 1. Experimental strategy, design and settings
- 2. Experimental Results
- 3. Future Plans.

MOTIVATION AND GOAL

- Use of a discrete model can be justified only when its solution converges to the solution of the original system as the resolution is refined.
- When truncation is made in or near an energetically-active range of the spectrum, model physics must also be changed as the resolution changes.
- At present, there is no unified formulation of model physics that automatically provides such changes.
- Since the original system is 3D, a converging framework must be at least quasi-3D.

MODEL RESOLUTION

HIERARCHY of QUASI-3D CRM AND MMF

3D CRM

Quasi-3D CRM

HIERARCHY of QUASI-3D CRM AND MMF

Quasi-3D CRM : A prerequisite to a quasi-3D MMF

A 4D estimation/prediction problem

with a highly-anisotropic multi-resolution horizontal grid with singular points ...

A treasure house of computational problems !

PATIENCE !

" If I have ever made valuable discoveries, it has been owing to patient attention, than to any other talent." Sir Isaac Newton

CLOUD-RESOLVING ANELASTIC MODEL BASED ON THE 3D VORTICITY EQUATION

 $u = \int^{\frac{1}{2}} \frac{\partial w}{\partial x}$

∂x $\int \frac{\partial w}{\partial x} + \rho_0 \eta$

 $\left(\frac{\partial w}{\partial x} + \rho_0 \eta\right)$

l

Prediction of scalar variables

Water substances (and tracer):

Potential temperature:

$$
\rho_0 \frac{\partial q_x}{\partial t} = -\left[\frac{\partial}{\partial x} (\rho_0 u q_x) + \frac{\partial}{\partial y} (\rho_0 v q_x) + \frac{\partial}{\partial z} (\rho_0 w q_x) \right] + S_{qx}
$$
\n
$$
\rho_0 \frac{\partial \theta}{\partial t} = -\left[\frac{\partial}{\partial x} (\rho_0 u \theta) + \frac{\partial}{\partial y} (\rho_0 v \theta) + \frac{\partial}{\partial z} (\rho_0 w \theta) \right] + S_{\theta}
$$

Prediction of vorticity components

Horizontal components:

$$
\rho_{o} \frac{\partial \xi}{\partial t} = -\left[\frac{\partial}{\partial x}(\rho_{o}u\xi) + \frac{\partial}{\partial y}(\rho_{o}v\xi) + \frac{\partial}{\partial z}(\rho_{o}w\xi)\right] + \rho_{o}\xi\frac{\partial u}{\partial x} + \rho_{o}\eta\frac{\partial u}{\partial y} + \rho_{o}\xi\frac{\partial u}{\partial z} + S_{\xi}
$$
\n
$$
\rho_{o} \frac{\partial \eta}{\partial t} = -\left[\frac{\partial}{\partial x}(\rho_{o}u\eta) + \frac{\partial}{\partial y}(\rho_{o}v\eta) + \frac{\partial}{\partial z}(\rho_{o}w\eta)\right] + \rho_{o}\eta\frac{\partial v}{\partial y} + \rho_{o}\xi\frac{\partial v}{\partial x} + \rho_{o}\xi\frac{\partial v}{\partial z} + S_{\eta}
$$
\n
$$
(\rho_{o}\zeta)_{z} = -\int_{z_{T}}^{z} \left[\frac{\partial(\rho_{o}\xi)}{\partial x} + \frac{\partial(\rho_{o}\eta)}{\partial y}\right] dz + (\rho_{o}\zeta)_{z_{T}}
$$

Vertical component:

Determination of wind components

From anelasticity:

$$
\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right) w + \frac{\partial}{\partial z} \left[\frac{1}{\rho_0} \frac{\partial}{\partial z} (\rho_0 w)\right] = -\rho_0 \frac{\partial \eta}{\partial x} + \rho_0 \frac{\partial \xi}{\partial y}
$$

$$
\int^z \left(\frac{\partial w}{\partial x} + \rho_0 \eta \right) dz + u_\tau(x, y, t) \qquad v = \int_{-\infty}^z \left(\frac{\partial w}{\partial x} - \rho_0 \xi \right) dz + v_\tau(x, y, t)
$$

 $\int_{z_{\text{T}}}^{z} \left(\frac{\partial \text{W}}{\partial y} - \rho_0 \xi \right) dz + v_{\text{T}}(x, y, t)$

From the definition of vorticity: $u = \int_{z_T}^{z} \left(\frac{\partial w}{\partial x} + \rho_0 \eta \right) dz + u_T(x, y, t)$ $v = \int_{z_T}^{z} \left(\frac{\partial w}{\partial y} - \rho_0 \xi \right) dz + u_T(x, y, t)$

A Problem with the Highly Anisotropic Grid :

Error imbalance between explicitly -resolved tangential derivative and statistically-estimated normal derivative.

The degree of imbalance is scale-dependent.

MULTI-SCALE REPRESENTATION OF VARIABLES

DETERMINATION OF 3D STRUCTURES

Determined by interpolatin of GCM grid-point values

(Currently, this field is prescribed.)

Along grid-point array: 1D Raynolds averaging of q *

Normal to grid-point array: By statistical Identification of cloud regime

q q' q"

Along grid-point array: q^* - q'

Normal to grid-point array: With a parameterization based on isotropy or inferred anisotropy

In CRMs with a horizontal resolutions of ≥ 1 km,

parameterization problems still exist *even for deep convection* parameterization problems still exist *even for deep convection*

(due to the existence of the internal structure and ajacent small-scale processes such as entrainment, detrainment, convection in anvils, convective downdrafts, triggering due to surface inhomogeneities, etc.)

Additional parameterization problem exists in MMFs, which have cloud-scale resolution only in limited directions

Emphasis of the Development Work up to Now

- *Cloud regimes have longer spatial and temporal scales than individual clouds.*
- *Cloud regimes can be statistically inferred through regression analysis of past data at the intersection and neighboring points.*

Technical Problems in Quasi3D Advection

Computational problems

- 1. Global stability with 2-dimensional uniform current
- 2. Local stability with 3-dimensional non-uniform current
- 3. Control of singularity at intersections
- 4. Control of spurious trend
- 5. Control of noise
- (6. Conservation)

Regression analysis of past data to identify cloud regimes

Solving 3D Elliptic Equation using the Quasi-3D Network

In anelastic models, 3D elliptic equation must be solved.

In momentum-equation models : for pressure

In vorticity-equation models : for vertical velocity

In our model, the elliptic equation for w is converted to a parabolic equation whose equilibrium \bullet solution is the solution of the elliptic equation (mimicing the relaxation method).

$$
\mu \frac{\partial w}{\partial t} = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right) w + \frac{\partial}{\partial z} \left[\frac{1}{\rho_0} \frac{\partial}{\partial z} \left(\rho_0 w\right)\right] + \rho_0 \left(\frac{\partial \eta}{\partial x} - \frac{\partial \xi}{\partial y}\right)
$$

The second-order finite difference term is estimated based on ghost point values given by

w': Statistically estimated as in advection

w": Currently assumed to be zero

No local deformation of the horizontal vorticity vector for cloud scale. \bullet

Problem with the Stretching and Twisting Terms in the Vorticity Equation

$$
\rho_0 \frac{\partial \xi}{\partial t} = -\left[\frac{\partial}{\partial x} (\rho_0 u \xi) + \frac{\partial}{\partial y} (\rho_0 v \xi) + \frac{\partial}{\partial z} (\rho_0 w \xi) \right] + \rho_0 \xi \frac{\partial u}{\partial x} + \rho_0 \eta \frac{\partial u}{\partial y} + \rho_0 \xi \frac{\partial u}{\partial z} + S_{\xi}
$$
\n
$$
\rho_0 \frac{\partial \eta}{\partial t} = -\left[\frac{\partial}{\partial x} (\rho_0 u \eta) + \frac{\partial}{\partial y} (\rho_0 v \eta) + \frac{\partial}{\partial z} (\rho_0 w \eta) \right] + \rho_0 \eta \frac{\partial v}{\partial y} + \rho_0 \xi \frac{\partial v}{\partial x} + \rho_0 \xi \frac{\partial v}{\partial z} + S_{\eta}
$$
\n
$$
\text{stretching} \quad \text{twisting}
$$

- These terms, which are responsible for enstrophy increase and energy cascade in 3D flow, are missing in 2D models.
- The Q3D CRM includes these terms.

Dilemma

- Fully for filtered scale (i.e., for cloud-system scale) vorticity;
- Partially for non-filtered scale (i.e., individual cloud scale) vorticity by neglecting the effect of local deformation on this scale.
- The enstrophy increase and energy cascade must be balanced by dissipation on average. So far we have not found a satisfactory formulation of the dissipation for the quasi-3D network.

Too strong dissipation : Inactive dynamics Too weak dissipation: Blowing up

Horizontal Advection of Precipitates in 3D CRM

∆x If $\Delta X < \frac{U}{V_t} Z_p \sim 10$ km, horizontal advection of precipitates must be explicitly formulated.

Lagrangian view: All precipitates in the column move into the next column before reaching ground. Eularian view: Air of the column is entirely replaced by the upwind-side precipitate-free air.

- If the precipitating system is poorly resolved (e.g., if ∆x ≥ 1 km) by an Eulerian grid, serious dispersion error may qppear.
- Consequently, a significant part of the precipitates can be left in the original column, which tends \bullet to evaporate as a result of drier-air inflow.

Thus in a CRM with 1 km $\leq \Delta x < \frac{U}{Vt}$ Z_p \sim 10 km, the area average of precipitate *due to convection tends to be under-predicted, and the humidity of the convective column tends to be over-predicted, eventually producing stratiform rain.*

These tendencies can be exaggerated in a Q3D CRM through

- Under-estimation of the precipitate mixing ratio at the up-wind side due to the use of statistics including non-precipitating cases.
- Overestimation of the normal component of velocity.

In our tets, generally the variance of water-vapor mixing ratio over the grid points tends to be over-predicted and those of rain/graupel mixing ratios tend to be under-predicted.

Unification of Dynamics between GCM and CRM

Two possibilities :

- (1) Use of the fully-compressible nonhydrostatic system of equations for both.
- (2) Use of a generalized anelastic system of equations for both, with compressibility included only for quasi-static motions.

This system automsticlly becomes the primitive equation model for large scales.