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Eddy-diffusivity/Mass-flux PBL parameterization: Dry convection and stratocumulus cases

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Motivation

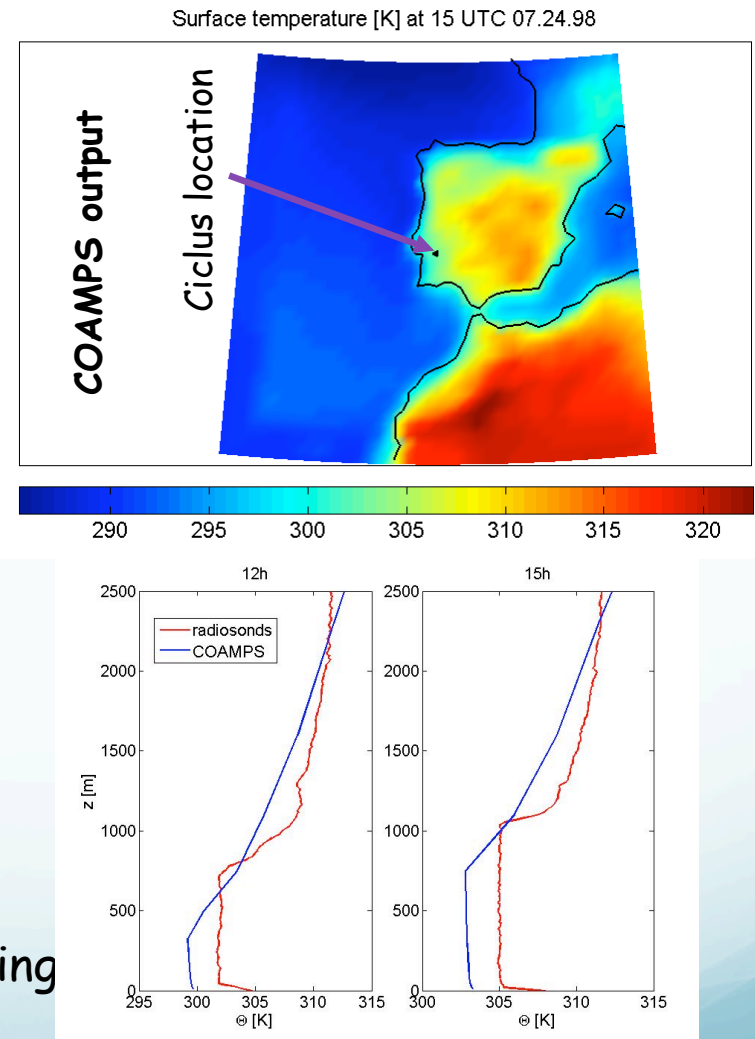
To improve PBL parameterizations in
weather forecast and climate models

- Dry convective boundary layer
- Shallow cumulus
- Stratocumulus
- Deep convection

Numerical models still have problems
simulating dry boundary layers

Realistic simulation of dry convection
is an essential stepping stone for tackling
cloudy boundary layers

Coupled Ocean-Atmosphere Mesoscale Prediction System (COAMPS) simulations





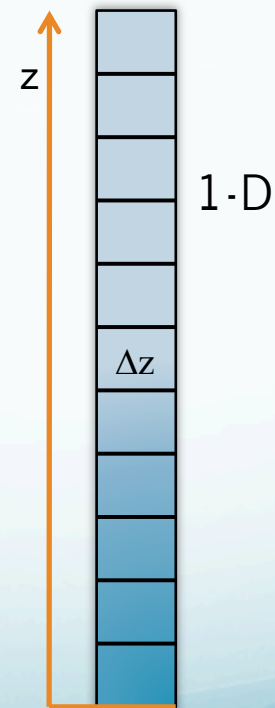
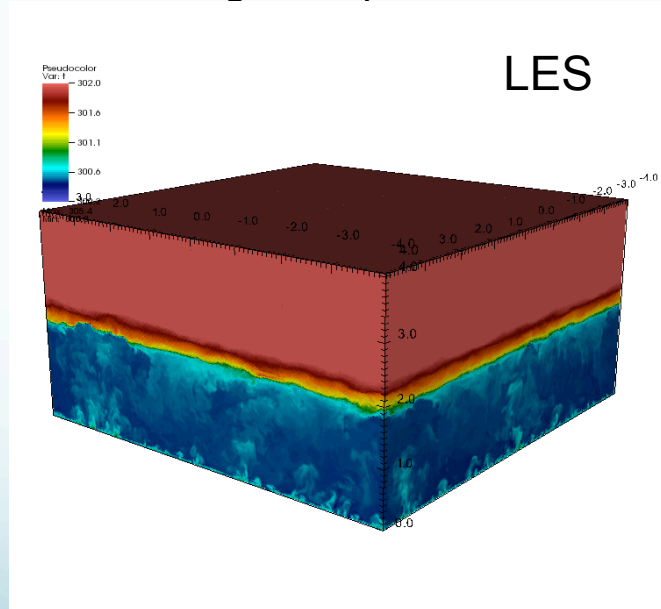
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Strategy

- more coherent and effective parameterizations
- Eddy-diffusivity/Mass-flux (EDMF) approach
- TKE prognostic equation
- Surface stability scaling
- PDF based cloud scheme

Tools: 1-D and Large Eddy Simulation (LES) models





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Dry convective PBL: free convection limit

Scalar prognostic equation:
$$\frac{\partial \bar{\phi}}{\partial t} = -\frac{\partial \overline{w'\phi'}}{\partial z} + F_\phi$$

EDMF concept :

(Siebesma and Teixeira, 2000)

$$\overline{w'\phi'} \equiv -K_\phi \frac{\partial \bar{\phi}}{\partial z} + M(\phi_u - \bar{\phi})$$

Mass-flux term:

$$M = \rho \sigma w_u \quad \text{Unknown: } \sigma, w_u, \phi_u, \varepsilon$$

Assumption: $\sigma = 0.1$ – we simulate the strongest 10% of the w distribution

$$w_u \frac{dw_u}{dz} = -\varepsilon \Delta_1 w_u^2 + \Delta_2 F_{w,u}$$

$$\frac{\partial \phi_u}{\partial z} = -\varepsilon (\phi_u - \bar{\phi})$$

$$\varepsilon = \alpha \frac{1}{l}$$

Entrainment coefficient – exchange of properties between updrafts and environment

eddy-diffusivity mixing (small eddies)

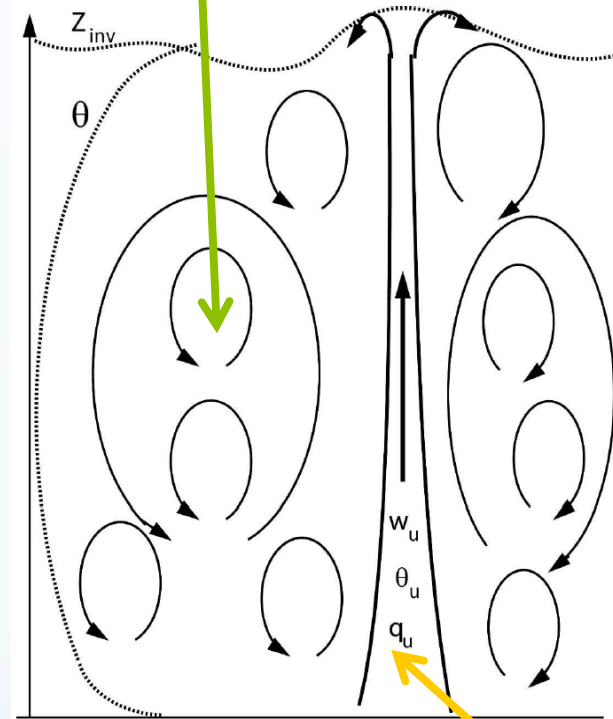


Fig. source: Siebesma et al. (2007)

mass-flux transport (strong buoyant updrafts)



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Eddy-diffusivity term:

$$K_{\phi,e} = C_k l \sqrt{e}$$

$$\overline{w'\phi'} \cong -K_\phi \frac{\partial \bar{\phi}}{\partial z} + M(\phi_u - \bar{\phi})$$

$$l^{-1} = l_1^{-1} + l_2^{-1}$$

$$l_1 = \tau \sqrt{e} \quad - \text{Teixeira and Cheinet (2004)}$$

$$l_2 = kz f(z/L) \quad - \text{e.g. Nakanishi (2001)}$$

$$\tau = a \frac{z_{inv}}{w_*}$$

- Important parameters: l and τ
- Closure based on e , surface stability, surface sensible heat flux and inversion height
- Minimal number of ad-hoc parameters

$$\frac{\partial e}{\partial t} = -\frac{\partial}{\partial z} \left(-K_e \frac{\partial e}{\partial z} \right) + \frac{g}{\theta} \overline{w'\theta'_v} - D$$

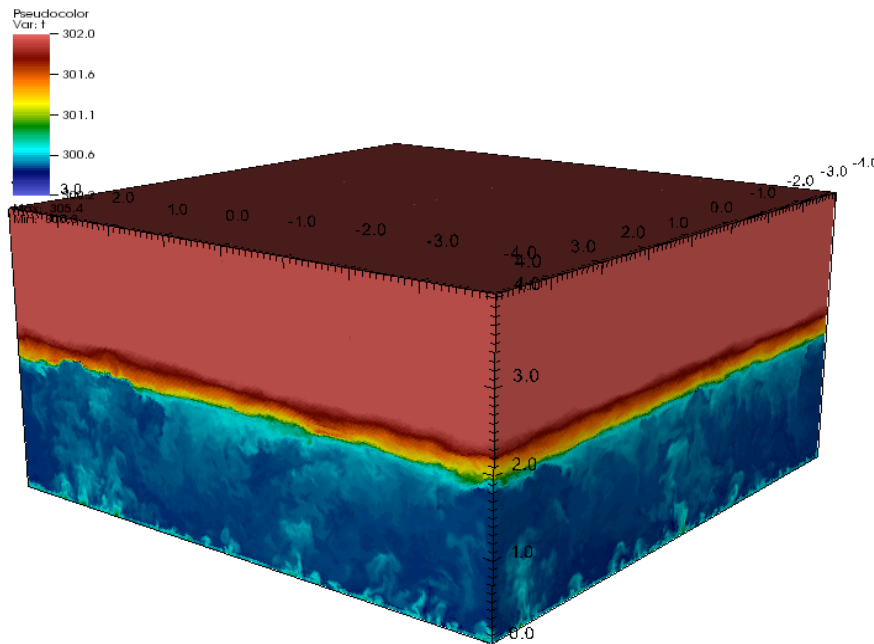
$$D = C_e e^{3/2} / l$$



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LES experiments: dry convective boundary layer with various surface heat fluxes



LES setup:

$$\Delta x = \Delta y = \Delta z = 20\text{m}$$

$$8 \times 8 \times 4(5) \text{ km}$$

$$\text{SHF}_s = (0.03, 0.06, 0.09, 0.12) [\text{Km/s}]$$

$$\text{EVP}_s = 2.5 \times 10^{-5} [\text{m/s}]$$

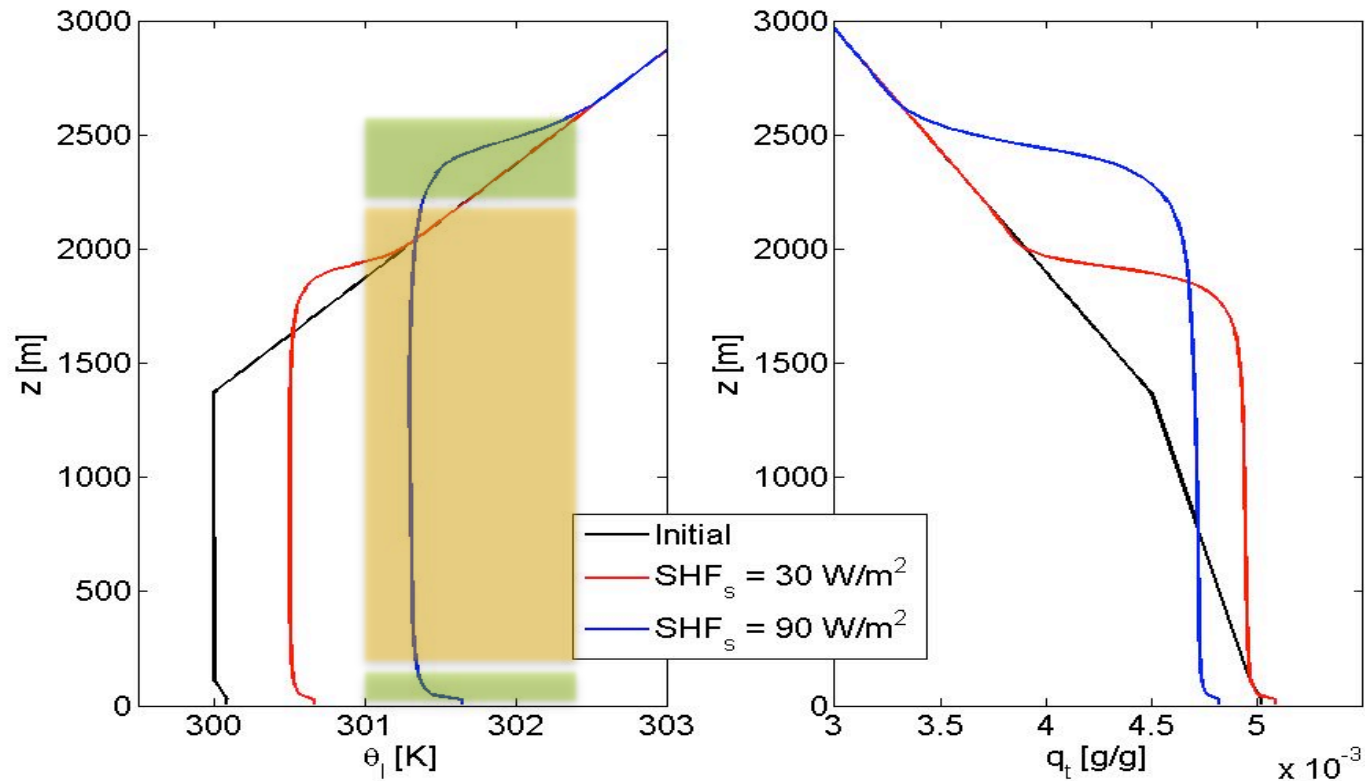
Initial θ and q_t profiles – next slide



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θ and q_t vertical profiles after 6 hours



Typical dry convective boundary layer structure:

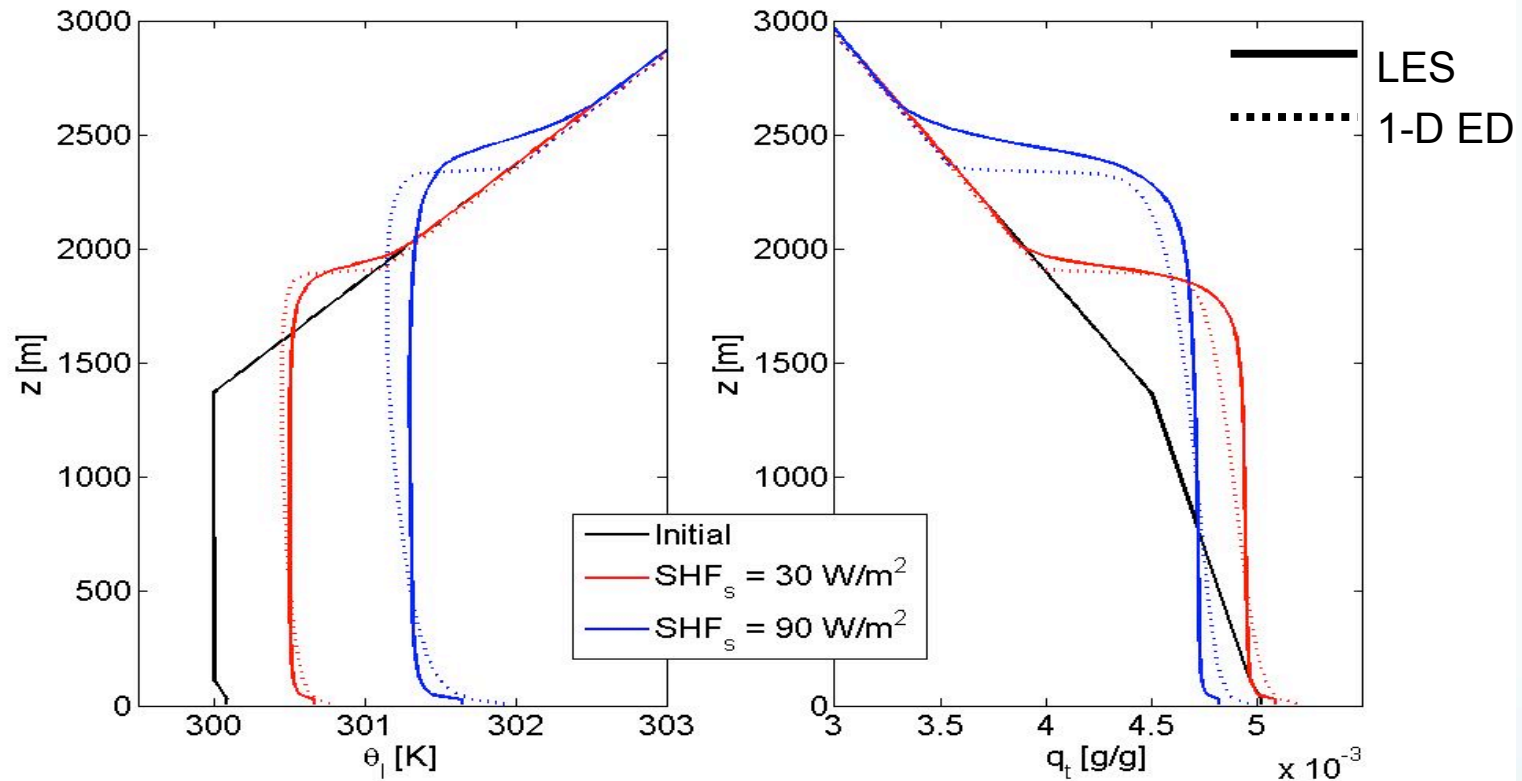
- unstable surface layer,
- well mixed layer, and
- capping inversion



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θ and q_t vertical profiles after 6 hours



Eddy-diffusivity simulations:

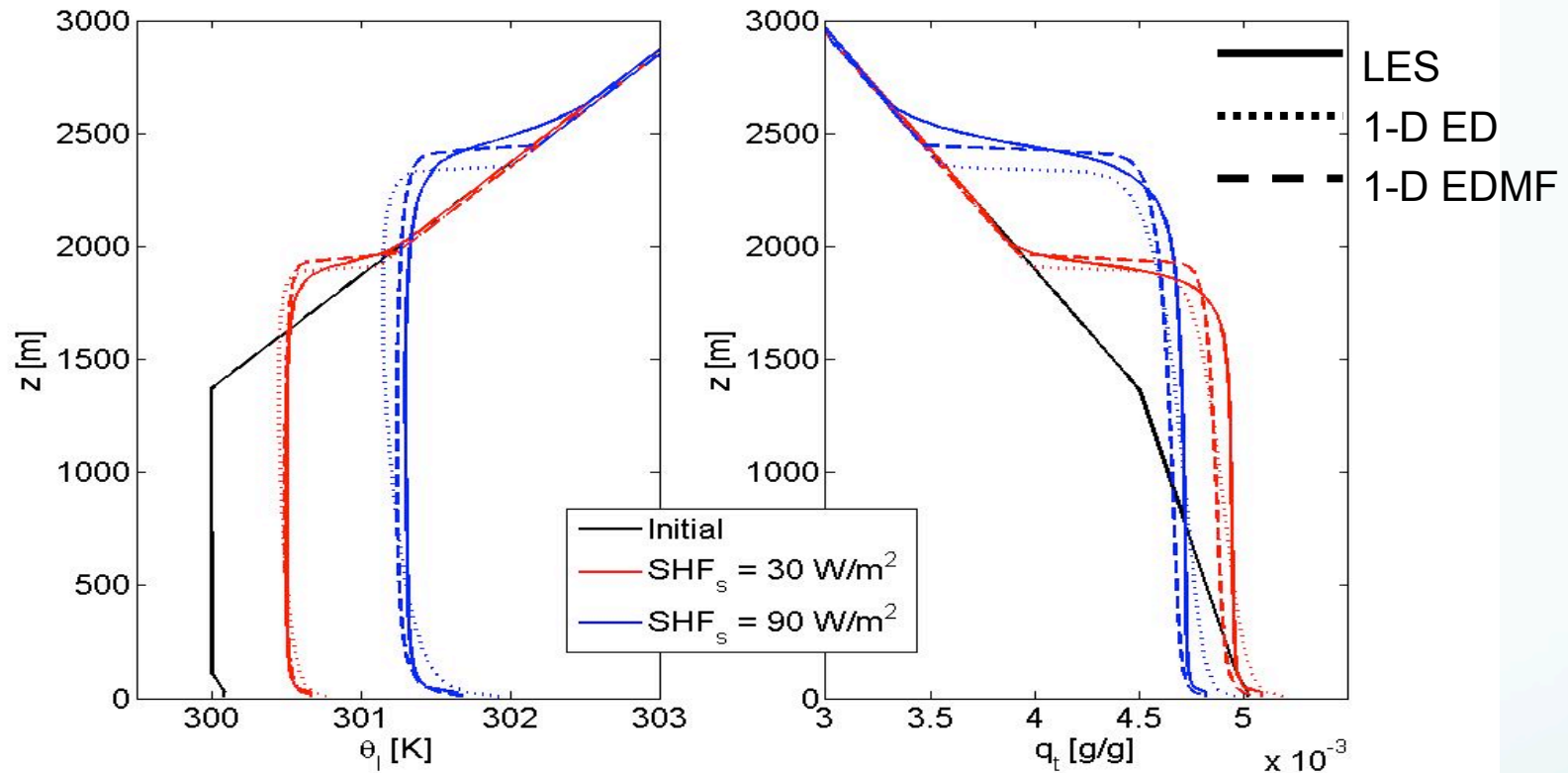
- the higher the surface heat flux the larger deviations from LES
- profiles not mixed well enough - unstable profile
- boundary layer too shallow
- surface profiles warmer than LES



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θ and q_t vertical profiles after 6 hours



Full EDMF simulations:

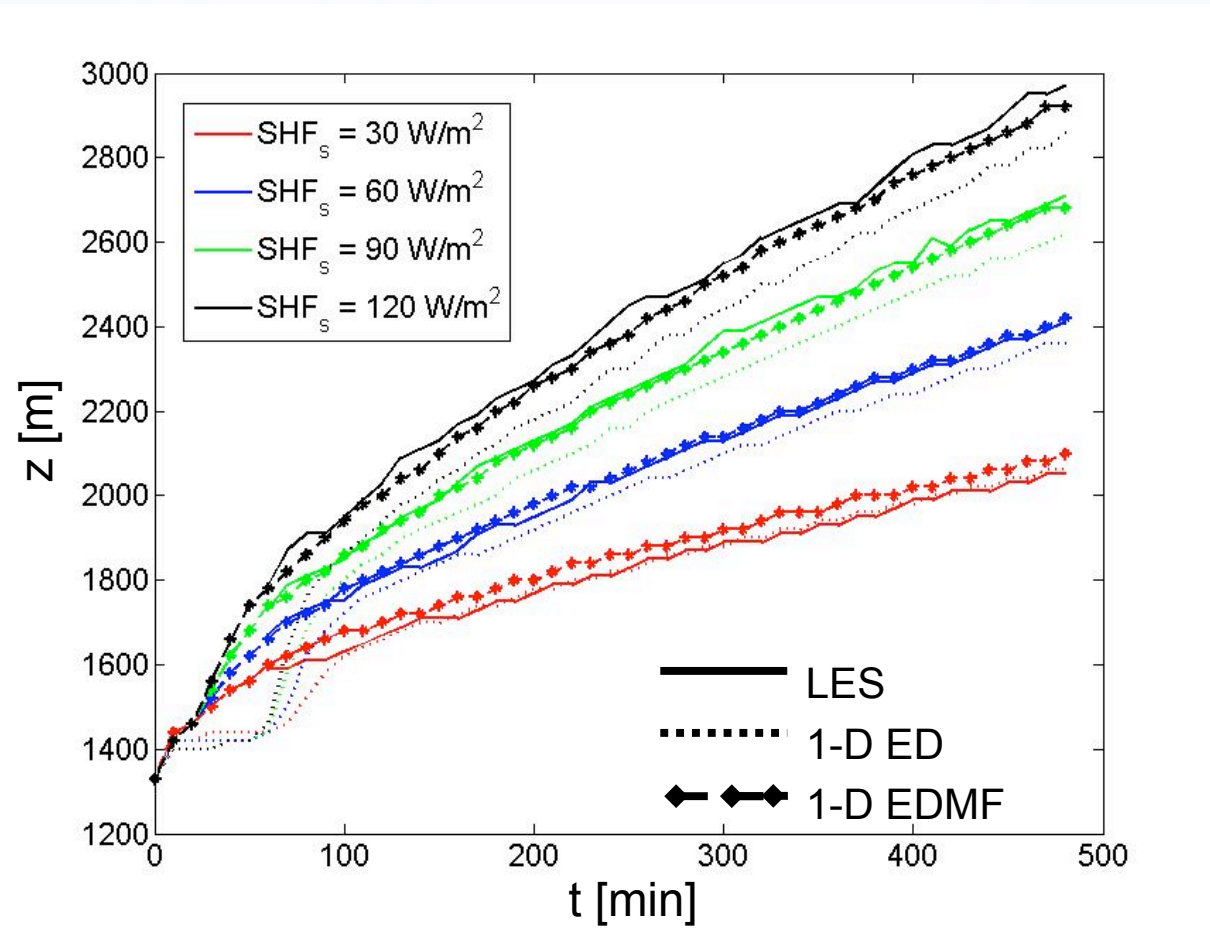
- surface layer more realistic
- neutral profile in the well-mixed layer
- larger entrainment leads to better inversion height
- inversion layer too sharp compared to LES



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Boundary layer height evolution



Inclusion of the mass-flux term improves
simulations of the boundary layer height for
a variety of surface heat fluxes



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Stratocumulus case

Moist physics

Prognostic equations for $\theta_L = \theta \left(1 - \frac{L_v}{T c_p} q_l\right)$ and $q_t = q_v + q_l$

PDF – based cloud parameterization (Cuijpers and Bechtold, 1995):

$$CC = 0.5 + 0.36 \arctan(1.55 Q_1)$$

$$q_l = \sigma_s \begin{cases} e^{1.2 Q_1 - 1} & Q_1 < 0 \\ e^{-1} + 0.66 Q_1 + 0.086 Q_1^2 & Q_1 > 0 \end{cases}$$

$$Q_1 = \frac{s}{\sigma_s} \quad \begin{array}{l} s - \text{saturation deficit} \\ \sigma_s - \text{standard deviation of } s \end{array}$$

$$s = q_t - q_s(T) = \frac{1}{1 + \gamma} (q_t - q_s(T_l))$$

$$\sigma_s = \max \left[l_s \left(\left(a^2 \frac{\partial q_t}{\partial z} \right)^2 + \left(b^2 \frac{\partial \theta_L}{\partial z} \right)^2 - 2ab \frac{\partial \theta_L}{\partial z} \frac{\partial q_t}{\partial z} \right)^{1/2}, 1e - 6 \right]$$



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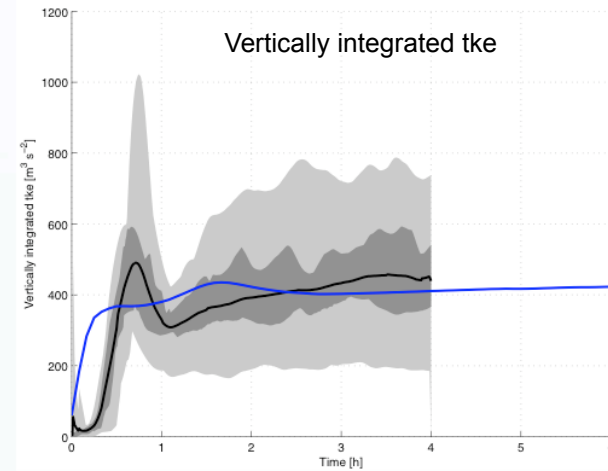
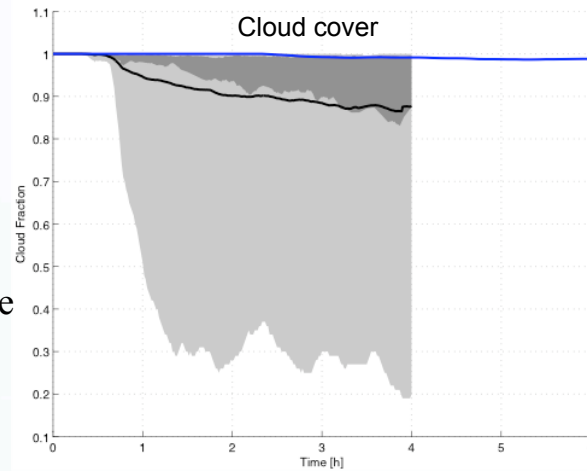
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DYCOMS stratocumulus case study - 1-D simulations vs. LES results (Stevens et al., 2005)

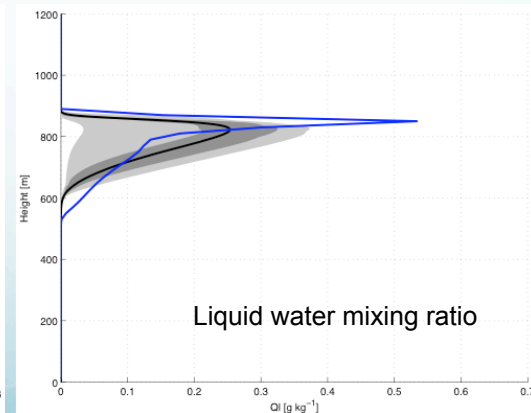
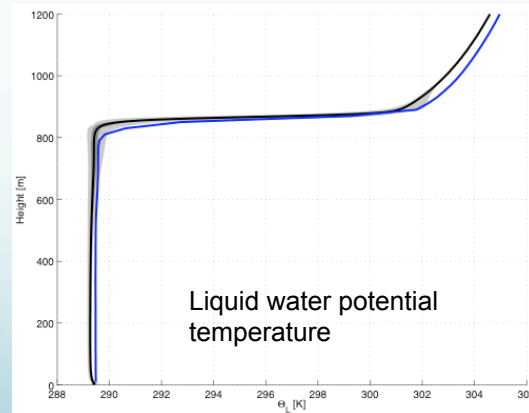
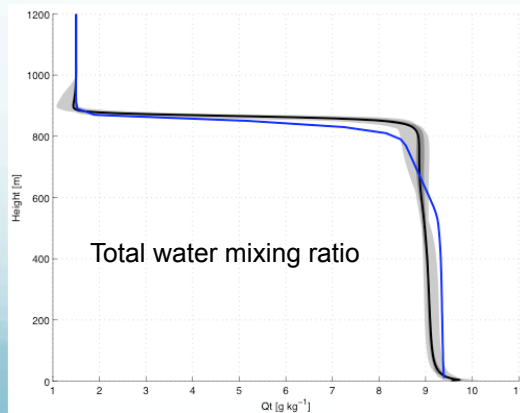
- Subtropical Pacific off-coast California
- Persistent cloud layer, homogeneous environmental conditions

Time series

— LES average
— 1-D sim.



Vertical profiles averaged between 3rd and 4th simulation hour:





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Conclusions and further plans

- EDMF improves 1D simulations of dry convective PBL's
- Stratocumulus simulations compare well with LES

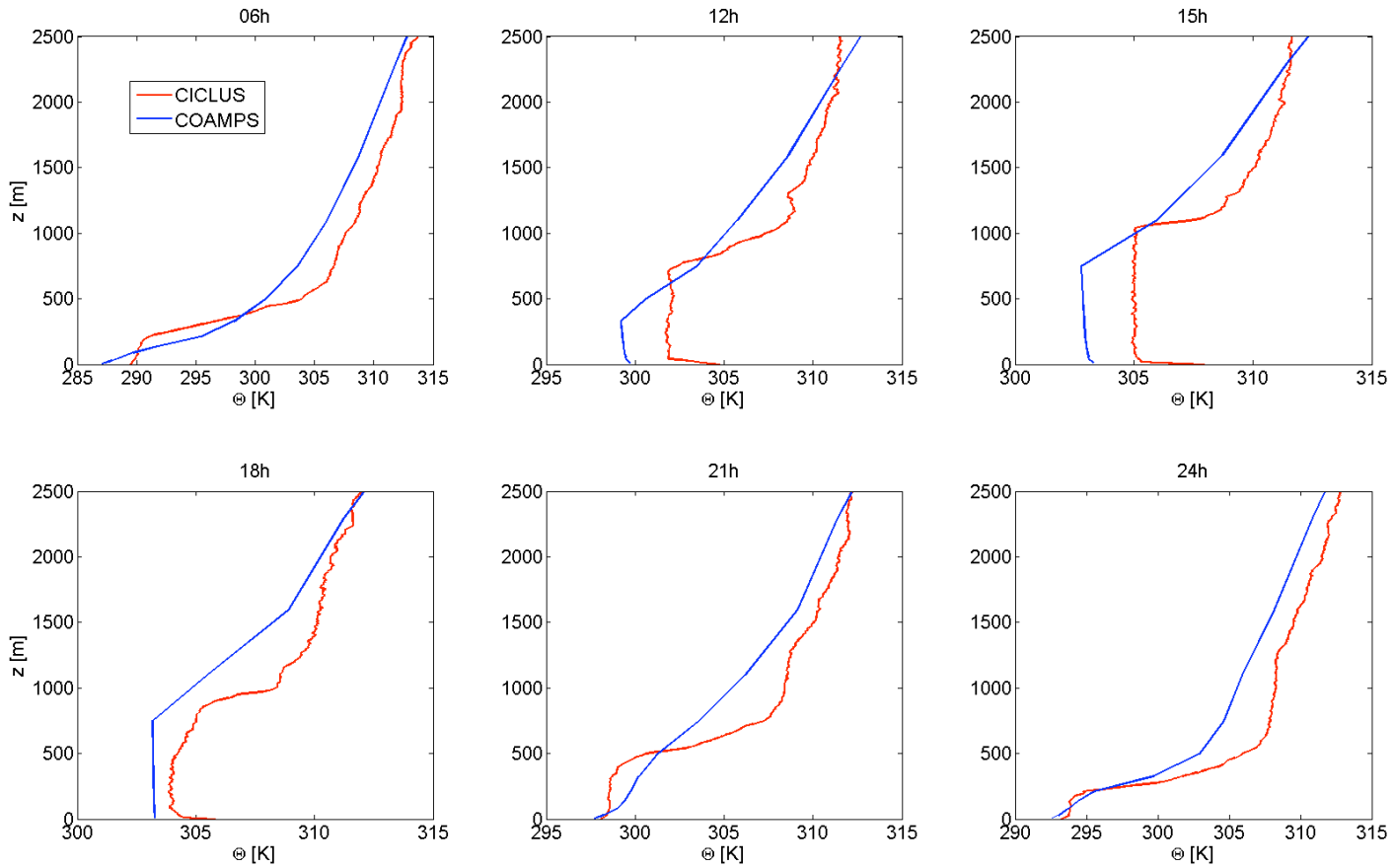
Plans:

- test EDMF parameterization in the COAMPS mesoscale model
- investigate simulation of Sc-to-Cu transition with EDMF



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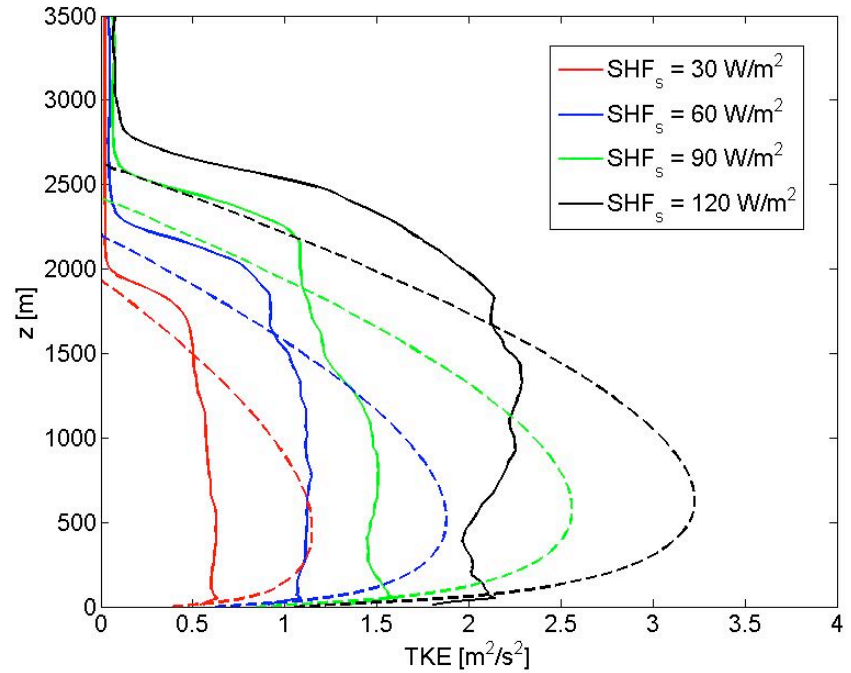
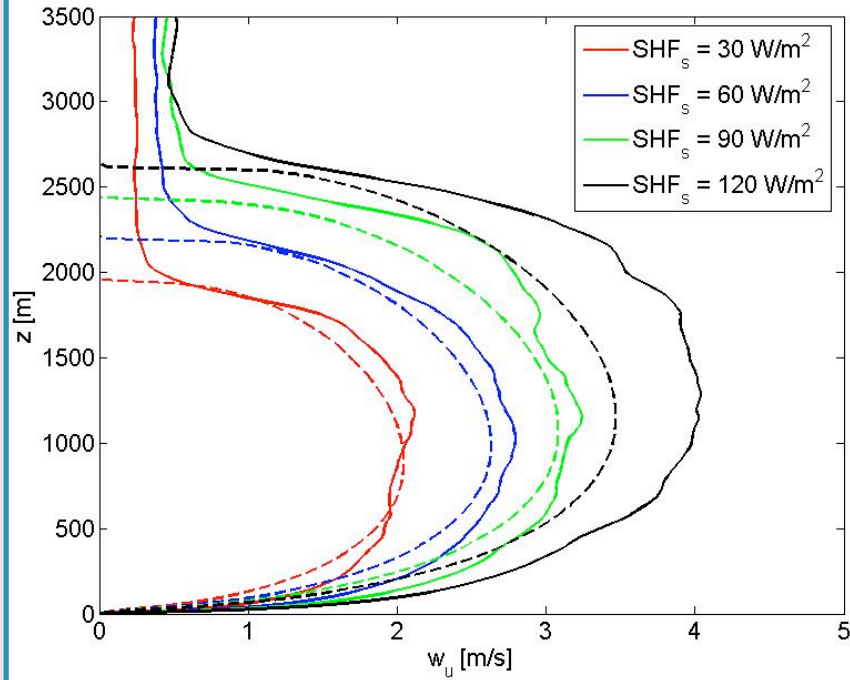
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$$\frac{\partial \phi_s}{\partial z} = -\epsilon(\phi_s - \bar{\phi})$$

$$\phi_{u,s} = \bar{\phi} + \beta \frac{\overline{w' \phi'^s}}{\sqrt{e_s}}, \quad \beta = 0.3$$

$$\frac{\partial w_u^2}{\partial z} = -\epsilon_w b_1 w_u^2 + b_2 F_{w,u}$$

$$F_{w,u} = g \left(\frac{\theta_{v,u}}{\theta_v} - 1 \right)$$





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DYCOMS initial conditions:

